

## Problem 1

- a) Yes, the state corresponds to a definite energy because the time-dependence is in the form compatible with stationary states:

$$e^{-\frac{ict}{\hbar}} \quad E = C$$

- b) for  $x < 0$   $\psi = 0 \Rightarrow V(x) = \infty$

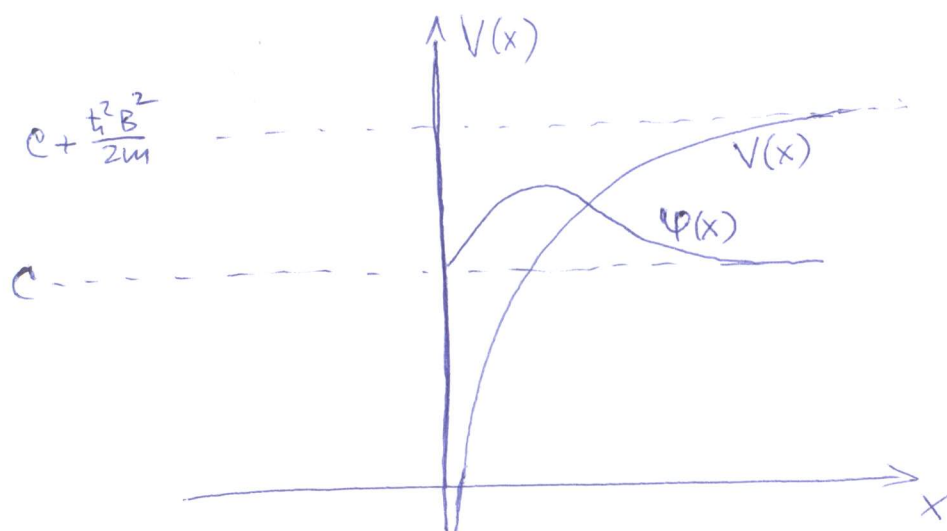
for  $x > 0$  we plug  $\psi$  into the Schrödinger equation. In fact, we can take just the spatial part of  $\psi(x,t)$  (let us call it  $\psi$ ) and plug it into the stationary Schrödinger equation:

$$\psi(x) = A x e^{-Bx} \quad (x > 0) \quad H\psi = C\psi$$

$$\frac{d\psi}{dx} = A e^{-Bx} - ABx e^{-Bx} \quad \frac{d^2\psi}{dx^2} = -AB e^{-Bx} - AB e^{-Bx} + AB^2 x e^{-Bx} =$$
$$= (-2B + B^2 x) A e^{-Bx} = \left(-\frac{2B}{x} + B^2\right) \psi$$

$$-\frac{\hbar^2}{2m} \left(-\frac{2B}{x} + B^2\right) \psi + V\psi = C\psi$$

$$V = C + \frac{\hbar^2}{2m} \left(B^2 - \frac{2B}{x}\right) = \left(C + \frac{\hbar^2 B^2}{2m}\right) - \frac{2\hbar^2 B}{2m} \frac{1}{x} \quad (x > 0)$$



Problem 2 We can rearrange  $V(x)$  in such a way that the linear term is eliminated:

$$V(x) = \frac{m\omega^2 x^2}{2} - bx = \frac{m\omega^2}{2} \left[ x^2 - \frac{2b}{m\omega^2} x \right] = \frac{m\omega^2}{2} \left[ \left( x - \frac{b}{m\omega^2} \right)^2 - \frac{b^2}{m^2\omega^4} \right] =$$
$$= \frac{m\omega^2}{2} (x - x_e)^2 + V_e$$

where  $x_e = \frac{b}{m\omega^2}$  and  $V_e = -\frac{b^2}{2m\omega^2}$

This potential is essentially the same as the ordinary harmonic oscillator potential. The only difference is that the energy scale is shifted by  $V_e$  and spatially the potential is shifted by  $x_e$  (which does not affect the energy). The solution of the Schrödinger equation is then obviously

$$E_n = V_e + \hbar\omega \left( n + \frac{1}{2} \right)$$

In particular the ground state energy is

$$E_0 = V_0 + \frac{\hbar\omega}{2} = -\frac{b^2}{2m\omega^2} + \frac{\hbar\omega}{2}$$

The ground state wave function is (see formula sheet)

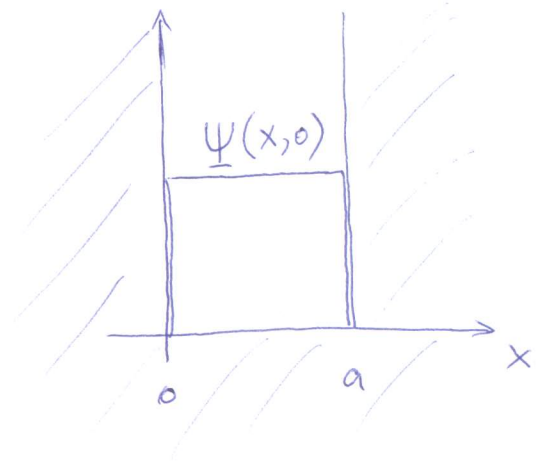
$$\psi(x) = \frac{\alpha^{1/4}}{\pi^{1/4}} e^{-\frac{\alpha(x-x_e)^2}{2}} \quad \alpha = \frac{m\omega}{\hbar}$$

### Problem 3

$$a) \quad 1 = \int_0^a |\Psi(x,0)|^2 dx = \int_0^a A^2 dx = A^2 a$$

$$A = \frac{1}{\sqrt{a}}$$

$$\text{Thus } \Psi(x,0) = \begin{cases} \frac{1}{\sqrt{a}}, & 0 < x < a \\ 0, & \text{otherwise} \end{cases}$$



$$b) \quad \Psi(x,t) = \sum_{n=1}^{\infty} c_n \psi_n(x) e^{-\frac{iE_n t}{\hbar}} \quad \text{where}$$

$$\psi_n(x) = \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi x}{a}\right)$$

$$E_n = \frac{n^2 \pi^2 \hbar^2}{2ma^2}$$

$$c_n = \int_0^a \psi_n^*(x) \Psi(x,0) dx = \int_0^a \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi x}{a}\right) \frac{1}{\sqrt{a}} dx =$$

$$= -\frac{\sqrt{2}}{a} \frac{a}{n\pi} \cos\left(\frac{n\pi x}{a}\right) \Big|_0^a = -\frac{\sqrt{2}}{\pi n} \underbrace{[\cos(n\pi) - 1]}_{(-1)^n} = \begin{cases} 0, & n \text{ is even} \\ \frac{\sqrt{8}}{\pi n}, & n \text{ is odd} \end{cases}$$

so

$$\Psi(x,t) = \sum_{k=0}^{\infty} \frac{4}{\pi \sqrt{a} (2k+1)} \sin\left(\frac{(2k+1)\pi x}{a}\right) e^{-i \frac{(2k+1)^2 \pi^2 \hbar}{2ma^2} t}$$

Problem 4  $H = \frac{p^2}{2m} + a|x|$

We can try to minimize the energy

$$E = \langle H \rangle = \frac{\langle p^2 \rangle}{2m} + a \langle |x| \rangle$$

$$(\Delta p)^2 = \langle p^2 \rangle - \langle p \rangle^2 \implies \langle p^2 \rangle \geq (\Delta p)^2$$

So we can use  $(\Delta p)^2$  instead of  $\langle p^2 \rangle$  as a lower bound in the expression for  $E$

Now  $\min \langle |x| \rangle \approx \Delta x$

then 
$$E \geq \frac{(\Delta p)^2}{2m} + a \Delta x$$

From the uncertainty principle we know  $\Delta p \Delta x \geq \frac{\hbar}{2}$  or  $\Delta x \approx \frac{\hbar}{2\Delta p}$

Hence, 
$$E \geq \frac{(\Delta p)^2}{2m} + \frac{a\hbar}{2\Delta p}$$

The minimum of the above expression is achieved when

$$\frac{\partial E}{\partial \Delta p} = 0 \implies \frac{\Delta p}{m} - \frac{a\hbar}{2(\Delta p)^2} = 0 \implies (\Delta p)^3 = \frac{a\hbar m}{2}$$

$$\Delta p = \left(\frac{a\hbar m}{2}\right)^{1/3}$$

$$E \geq \left(\frac{a\hbar m}{2}\right)^{2/3} \frac{1}{2m} + \frac{a\hbar}{2} \left(\frac{2}{a\hbar m}\right)^{1/3} = \frac{a^{2/3} \hbar^{2/3}}{m^{1/3}} \left(\frac{1}{2^{5/3}} + \frac{1}{2^{2/3}}\right) =$$

$$= \frac{3}{2^{5/3}} \frac{a^{2/3} \hbar^{2/3}}{m^{1/3}} = \frac{3}{2m} \left(\frac{a\hbar m}{2}\right)^{2/3}$$