

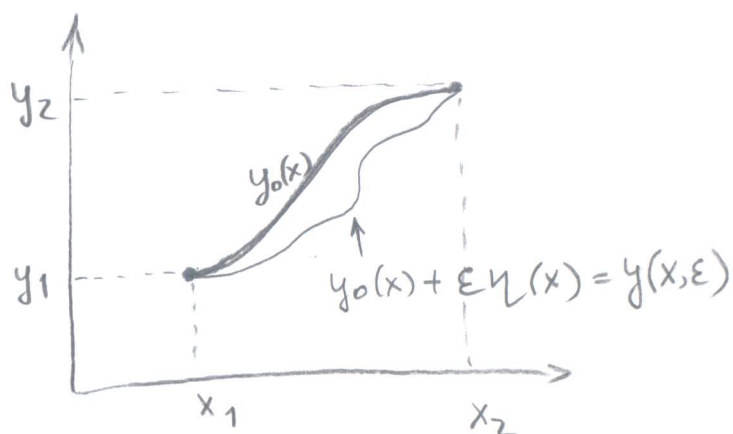
# General discussion of variational principles

Let us consider the following one-dimensional problem. Given an integrable function  $F(y, y', x)$  defined on a path  $y = y(x)$  between two values  $x_1$  and  $x_2$ , where  $y' = \frac{dy}{dx}$ , we wish to find a particular path such as the line integral

$$I = \int_{x_1}^{x_2} F(y, y', x) dx$$

takes an extremum value (e.g.  $I$  has a stationary value relative to path differing infinitesimally from the  $y(x)$  which realizes the extremum)

Our problem is transformed into an elementary extremum value problem by covering the ensemble of all physically meaningful paths by a parametric representation:



$$y(x, \epsilon) = y_0(x) + \epsilon \eta(x)$$

where  $\epsilon$  is a parameter constant for every path,  $\eta(x)$  is an arbitrary differentiable function that vanishes at  $x_1$  and  $x_2$ :  $\eta(x_1) = \eta(x_2) = 0$ .

The desired curve is given by  $y_0(x) = y(x, 0)$

For any such parametric family of curves  $I$  is also a function of  $\epsilon$

$$I(\epsilon) = \int_{x_1}^{x_2} F(y(x, \epsilon), y'(x, \epsilon), x) dx$$

The condition for an extremum value of integral  $I$  is then

$$\left. \frac{dI}{d\varepsilon} \right|_{\varepsilon=0} = 0$$

The differentiation under the integral symbol (allowed if  $F$  is continuously differentiable with respect to  $\varepsilon$ ) yields

$$\frac{dI}{d\varepsilon} = \int_{x_1}^{x_2} \left( \frac{\partial F}{\partial y} \frac{\partial y}{\partial \varepsilon} + \frac{\partial F}{\partial y'} \frac{\partial y'}{\partial \varepsilon} + \frac{\partial F}{\partial x} \frac{\partial x}{\partial \varepsilon} \right) dx = \int_{x_1}^{x_2} \left( \frac{\partial F}{\partial y} \eta + \frac{\partial F}{\partial y'} \eta' \right) dx$$

The second term can be integrated by parts

$$\int_{x_1}^{x_2} \frac{\partial F}{\partial y'} \frac{\partial \eta}{\partial x} dx = \left[ \frac{\partial F}{\partial y'} \eta \right]_{x_1}^{x_2} - \int_{x_1}^{x_2} \left( \frac{d}{dx} \frac{\partial F}{\partial y'} \right) \eta dx$$

0, because  $\eta(x_1) = \eta(x_2) = 0$

Then the extremum condition becomes

$$\int_{x_1}^{x_2} \left( \frac{\partial F}{\partial y} - \frac{d}{dx} \frac{\partial F}{\partial y'} \right) \eta dx = 0$$

The "fundamental lemma" of the calculus of variations says if

$$\int_{x_1}^{x_2} G(x) \eta(x) dx = 0$$

for any arbitrary function  $\eta(x)$  continuous through the second derivative then  $G(x)$  must vanish. Therefore

We can write

$$\frac{\partial F}{\partial y} - \frac{d}{dx} \left( \frac{\partial F}{\partial y'} \right) = 0$$

← Euler-Lagrange equation

The solution of the Euler-Lagrange equation, a differential equation of second order, together with boundary conditions, yields the path sought.

Let us define the variation of a function  $y(x, \epsilon)$  as

$$\delta y = y(x, d\epsilon) - y(x, 0) = \left. \frac{\partial y}{\partial \epsilon} \right|_{\epsilon=0} \cdot d\epsilon \quad d\epsilon \rightarrow 0$$

Then the variational problem can be formulated as

$$\delta \int_{x_1}^{x_2} F(y(x), y'(x), x) dx = 0$$

Function  $F$  can also include constraints by means of Lagrange multipliers.

### Generalization to many independent variables and derivation of Lagrange's equations

A variation of the integral  $I$  in case of multiple independent variables  $y_i(x)$

$$\delta I = \delta \int_{x_1}^{x_2} F(y_1(x), \dots, y_n(x); y_1'(x), \dots, y_n'(x), x) dx$$

is obtained by considering  $I$  as a function of a parameter that labels sets of curves

$$y_1(x, \epsilon) = y_1(x, 0) + \epsilon \eta_1(x)$$

$\vdots$

$$y_n(x, \epsilon) = y_n(x, 0) + \epsilon \eta_n(x)$$

where  $\eta_1 \dots \eta_n$  are independent arbitrary functions that vanish at  $x_1$  and  $x_2$

$$\delta I = \frac{\partial I}{\partial \epsilon} d\epsilon = \int_{x_1}^{x_2} \sum_{i=1}^n \left( \frac{\partial F}{\partial y_i} \frac{\partial y_i}{\partial \epsilon} + \frac{\partial F}{\partial y_i'} \frac{\partial y_i'}{\partial \epsilon} d\epsilon \right) dx$$

Again we integrate by parts the second term

$$\int_{x_1}^{x_2} \frac{\partial F}{\partial y_i'} \frac{\partial^2 y_i}{\partial \epsilon \partial x} dx = \underbrace{\frac{\partial F}{\partial y_i'} \frac{\partial y_i}{\partial \epsilon}}_0 \Big|_{x_1}^{x_2} - \int_{x_1}^{x_2} \frac{\partial y_i}{\partial \epsilon} \frac{d}{dx} \left( \frac{\partial F}{\partial y_i'} \right) dx$$

Then

$$\delta I = \int_{x_1}^{x_2} \sum_{i=1}^n \left( \frac{\partial F}{\partial y_i} - \frac{d}{dx} \frac{\partial F}{\partial y_i'} \right) \delta y_i dx \quad \text{where } \delta y_i = \frac{\partial y_i}{\partial \epsilon} \Big|_{\epsilon=0} \cdot d\epsilon$$

By the extension of the "fundamental lemma" the condition of extrema requires that all coefficients by  $\delta y_i$  vanish:

$$\frac{\partial F}{\partial y_i} - \frac{d}{dx} \left( \frac{\partial F}{\partial y_i'} \right) = 0 \quad i = 1, \dots, n$$

If we now change notations and replace

$$F \rightarrow L \quad y_i \rightarrow q_i \quad y_i' \rightarrow \dot{q}_i \quad x \rightarrow t$$

then the principle of least action  $\delta I = 0$

$$\text{with } I = \int_{t_1}^{t_2} L(q_i, \dot{q}_i, t) dt$$

yields

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i} - \frac{\partial L}{\partial q_i} = 0 \quad i = 1, \dots, n$$

under assumption that  $q_i$  are all independent, which requires that the constraints be holonomic